



#### Thermomechanical and Interfacial Properties of Graphene Membranes

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# Membrane Materials or Structures

- Mechanical (solid) membranes: no bending rigidity
- Bio-membranes (lipid bilayers, fluid membranes, etc.)
- 2D crystal monolayers: graphene, h-BN, MoS2, etc.



#### **Nonlinear Continuum Mechanics of 2D Sheets**

2D-to-3D deformation gradient:

$$F_{iJ} = \frac{\partial x_i}{\partial X_J}$$



In-plane deformation: 2D Green-Lagrange strain tensor

Bending: 2D curvature tensor (strain gradient)

$$\mathbf{K}_{IJ} = n_i \frac{\partial F_{iI}}{\partial X_J} = n_i \frac{\partial^2 x_i}{\partial X_I \partial X_J} \implies \kappa_{\alpha\beta} = \frac{\partial^2 w}{\partial x_\alpha \partial x_\beta}$$
  
Strain energy: 
$$U = \int_A \Phi(\mathbf{E}, \mathbf{K}) dA$$
 Temperature effect?

Lu and Huang, Int. J. Applied Mechanics 1, 443-467 (2009).

# **Small deformation at T = 0 K (Statics)** <u>Linear elastic strain energy</u>: $\Phi(\varepsilon, \kappa) = \Phi_{\varepsilon}(\varepsilon) + \Phi_{\kappa}(\kappa)$

In-plane stretching: 
$$\Phi_{\varepsilon} = \frac{E}{2(1+\nu)} \left( \varepsilon_{\alpha\beta} \varepsilon_{\alpha\beta} + \frac{\nu}{1-\nu} \varepsilon_{\alpha\alpha} \varepsilon_{\beta\beta} \right)$$

Bending:  

$$\Phi_{\kappa} = \frac{1}{2} D (\kappa_1 + \kappa_2)^2 - D_G \kappa_1 \kappa_2$$

Moderately nonlinear kinematics:  $\varepsilon_{\alpha\beta} = \frac{1}{2} \left( \frac{\partial u_{\alpha}}{\partial x_{\beta}} + \frac{\partial u_{\beta}}{\partial x_{\alpha}} + \frac{\partial w}{\partial x_{\alpha}} \frac{\partial w}{\partial x_{\beta}} \right)$ 

The basic elastic properties  $(E, v, D, \text{ and } D_G)$  can be determined by DFT or molecular mechanics (statics) calculations. <u>They are</u> <u>NOT directly related to each other!</u>

#### Elastic properties of graphene at T = 0 K



Method	E (N/m)	V	E/(1-v)	D (eV)
DFT	345	0.149	406	1.5
MM (REBO-2)	243	0.397	403	1.4
MM (AIREBO)	277	0.366	437	1.0

### T > 0 K: a hybrid approach



#### Harmonic analysis of thermal fluctuation

$$w(\mathbf{r}) = \sum_{k} \hat{w}(\mathbf{q}_{k}) e^{i\mathbf{q}_{k}\cdot\mathbf{r}} \qquad U_{b} = \frac{DL_{0}^{2}}{2} \sum_{k} q_{k}^{4} \left[ \hat{w}_{\text{Re}}^{2}(\mathbf{q}_{k}) + \hat{w}_{\text{Im}}^{2}(\mathbf{q}_{k}) \right]$$

$$Z_{b} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} \exp\left(-\frac{U_{b}}{k_{B}T}\right) d\hat{w}_{\text{Re}}(\mathbf{q}_{1}) d\hat{w}_{\text{Im}}(\mathbf{q}_{1}) \cdots = \prod_{k(\mathbf{q}_{k} \cdot \mathbf{e}_{y} \ge 0)} \left(\frac{\pi k_{B}T}{DL_{0}^{2}q_{k}^{4}}\right)$$

$$\left\langle \left| \hat{w}(\mathbf{q}_{k}) \right|^{2} \right\rangle = \frac{1}{Z_{b}} \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} \left| \hat{w}(\mathbf{q}_{k}) \right|^{2} \exp\left(-\frac{U_{b}}{k_{B}T}\right) d\hat{w}_{\text{Re}}(\mathbf{q}_{1}) d\hat{w}_{\text{Im}}(\mathbf{q}_{1}) \cdots = \frac{k_{B}T}{DL_{0}^{2}q_{k}^{4}}$$

$$\left\langle h^{2} \right\rangle = \sum \left\langle \left| \hat{w}(\mathbf{q}_{1}) \right|^{2} \right\rangle = \frac{k_{B}T}{k_{B}T} \sum \alpha^{-4}$$

$$\langle h^2 \rangle = \sum_k \left\langle \left| \hat{w}(\mathbf{q}_k) \right|^2 \right\rangle = \frac{\kappa_B r}{DL_0^2} \sum_k q_k^-$$

$$\overline{h} = \sqrt{\left\langle h^2 \right\rangle} = L_0 \sqrt{\frac{\gamma_n k_B T}{16\pi^4 D}}$$

Basically a linear plate model, independent of in-plane stiffness or Gaussian curvature.

# Effect of pre-tension (still harmonic) $U = U_b + U_s = E^* \varepsilon_0^2 L_0^2 + L_0^2 \sum_{k(\mathbf{q}_k \cdot \mathbf{e}_v \ge 0)} (Dq_k^4 + E^* \varepsilon_0 q_k^2) (\hat{w}_{\text{Re}}^2(\mathbf{q}_k) + \hat{w}_{\text{Im}}^2(\mathbf{q}_k))$

$$\left\langle h^{2} \right\rangle = \frac{k_{B}T}{DL_{0}^{2}} \sum_{k} q_{k}^{-4} \left( 1 + \frac{E^{*}\varepsilon_{0}}{Dq_{k}^{2}} \right)^{-1}$$
$$\approx \frac{k_{B}T}{4\pi E^{*}\varepsilon_{0}} \ln \left( 1 + \frac{E^{*}\varepsilon_{0}L_{0}^{2}}{4\pi^{2}D} \right)$$

A small pre-tension can solution considerably suppress thermal fluctuation, resulting in a different scaling behavior.

$$l = \sqrt{\frac{D}{E^* \varepsilon_0}} = \frac{0.024 \text{nm}}{\sqrt{\varepsilon_0}}$$



Gao and Huang, JMPS, in press.

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# **Nonlinear thermoelasticity**

Partition function:

$$Z(\varepsilon_0, T) = \exp\left(-\frac{E^*\varepsilon_0^2 L_0^2}{k_B T}\right) \prod_{k(\mathbf{q}_k \cdot \mathbf{e}_y \ge 0)} \left[ \left(1 + \frac{E^*\varepsilon_0}{Dq_k^2}\right)^{-1} \left(\frac{\pi k_B T}{DL_0^2 q_k^4}\right) \right]$$

T1----

Helmholtz free energy:  $A(\varepsilon_0, T) = -k_B T \ln Z(\varepsilon_0, T)$ 

Stress: 
$$\sigma(\varepsilon_0, T) = \frac{1}{2L_0^2} \left( \frac{\partial A}{\partial \varepsilon_0} \right)_T = E^* \varepsilon_0 + \widetilde{\sigma}(\varepsilon_0, T)$$
 induced  
Tangent modulus:  $\widetilde{E}^* = \frac{\partial \sigma}{\partial \varepsilon_0} \approx E^* - \frac{E^* k_B T}{16\pi D} \left( \frac{4\pi^2 D}{E^* L_0^2} + \varepsilon_0 \right)^{-1}$   
Thermal softening and strain stiffening; size dependent?

# **Thermal expansion**



- Anharmonic in-plane oscillations result in a positive thermal expansion, nearly independent of temperature (up to 1000 K).
- Out-of-plane fluctuation leads to in-plane contraction (negative thermal expansion) – anharmonic effects TBD



# MD Simulations (T > 0 K)

- Free-standing graphene in NPT ensemble (zero stress)
- Constrained graphene in NVT ensemble (zero strain)
- Biaxially strained graphene in NVT ensemble



- > Thermal rippling
- Thermal expansion/contraction
- Thermal stress
- Temperature dependent mechanical properties

# **Thermal Rippling**



Harmonic approximation:  $\overline{h} = \sqrt{\langle h^2 \rangle} = L_0 \sqrt{\frac{\gamma_n k_B T}{16\pi^4 D}}$ 

Significant anharmonic effects due to coupling between bending and stretching ( $\zeta < 1$ ), similar to biomembranes!

Gao and Huang, JMPS, in press.

# **NPT: Thermal Expansion/Contraction**



- Negative thermal expansion at low T, and positive at high T.
- Thermal expansion/contraction is size dependent!
- By suppressing out-of-plane fluctuations, 2D simulations predict a constant positive CTE (size-independent).

Gao and Huang, JMPS, in press.

# **NVT: Thermal Stress at Zero Strain**



- As expected, negative thermal expansion leads to tensile stress at low T, and the opposite is true at high T.
- However, thermal rippling differs under NPT and NVT.

# **NVT: biaxially strained graphene**



<u>Nonlinear elasticity due to two effects</u>:(1) intrinsic strain softening (large strain behavior);

(2) strain stiffening due to thermal rippling (small strain behavior)



Gao and Huang, JMPS, in press.

## **Biaxial modulus at zero strain**



Due to the effect of thermal rippling, the elastic modulus becomes both temperature and size-dependent at zero strain; this effect however is largely beyond harmonic.

Gao and Huang, JMPS, in press.

# **Biaxial modulus at 1% strain**



- The effect of thermal rippling reduces as the strain increases;
- The tangent elastic modulus decreases linearly with temperature (almost harmonic), and becomes size independent for relatively large membranes.
   *Gao and Huang, JMPS, in press.*

# **Graphene on substrate**



 Assume that the in-plane dimension of graphene follows thermal expansion of the substrate (full strain transfer).



Gao and Huang, JMPS, in press.

# **Interfacial strain transfer**



Jiang, Huang, and Zhu, Adv. Funct. Mater., 2014.

### Interfacial properties: Adhesion and Friction



$$U_{vdW}(z) = -\Gamma_0 \left[ \frac{3}{2} \left( \frac{h_0}{z} \right)^3 - \frac{1}{2} \left( \frac{h_0}{z} \right)^9 \right]$$



Aitken and Huang, J. Appl. Phys. 107, 123531 (2010).



- Zero friction on a perfectly flat surface (assuming vdW).
- Friction strength depends on surface roughness and adhesion.
- Temperature effect?

# A blister test



- CVD grown graphene was transferred to a copper substrate.
- The graphene/photoresist composite film was pressurized with deionized water.
- Deflection profiles were measured by a full field interference method.
- Energy release rate was calculated as a function of delamination growth to obtain fracture resistance curves.
- The delamination path was confirmed by Raman spectroscopy.

Cao et al., Carbon 2014.





## **Delamination Resistance Curves**



Compare to other measurements:

- Graphene/SiO<sub>2</sub>: 0.2-0.45 J/m<sup>2</sup> (Scott Bunch)
- Graphene/Cu film: 0.72 J/m<sup>2</sup> (Yoon et al.)

Cao et al., Carbon 2014.

#### **Graphene/SiO<sub>2</sub>: effect of surface structures**



# **Graphene/SiO<sub>2</sub>: vdW-DFT**





The vdW adhesion energy is reduced by surface hydroxylation and further reduced by adsorption of water molecules.

Gao et al., submitted.

## **Graphene/SiO<sub>2</sub>: capillary forces?**



# Summary

- A statistical mechanics approach is employed to study thermal rippling and thermoelasticity of graphene membranes; the current analysis is limited by harmonic analysis.
- MD simulations show significant anharmonic effects at zero stress or zero strain, but nearly harmonic behavior when the membrane is subjected to pretension.
- Graphene on substrate: strain transfer (shear) and adhesion have been measured, but the underlying mechanisms (vdW or capillary) require further studies.